



First Results from the CARIBU Facility: Mass Measurements on the r -Process Path

J. Van Schelt,^{1,2} D. Lascar,^{3,1} G. Savard,^{1,2} J. A. Clark,¹ P. F. Bertone,¹ S. Caldwell,^{2,1} A. Chaudhuri,^{4,1} A. F. Levand,¹ G. Li,^{5,1} G. E. Morgan,⁴ R. Orford,⁵ R. E. Segel,^{3,1} K. S. Sharma,⁴ and M. G. Sternberg^{2,1}

¹Physics Division, Argonne National Laboratory, Argonne, Illinois 60439, USA

²Department of Physics, University of Chicago, Chicago, Illinois 60637, USA

³Department of Physics and Astronomy, Northwestern University, Evanston, Illinois 60208, USA

⁴Department of Physics and Astronomy, University of Manitoba, Winnipeg, Manitoba R3T 2N2, Canada

⁵Department of Physics, McGill University, Montréal, Québec H3A 2T8, Canada

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The Canadian Penning Trap mass spectrometer has made mass measurements of 33 neutron-rich nuclides provided by the new Californium Rare Isotope Breeder Upgrade facility at Argonne National Laboratory. The studied region includes the ^{132}Sn double shell closure and ranges in Z from In to Cs, with Sn isotopes measured out to $A = 135$, and the typical measurement precision is at the 100 ppb level or better. The region encompasses a possible major waiting point of the astrophysical r process, and the impact of the masses on the r process is shown through a series of simulations. These first-ever simulations with direct mass information on this waiting point show significant increases in waiting time at Sn and Sb in comparison with commonly used mass models, demonstrating the inadequacy of existing models for accurate r -process calculations.

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Introduction.—Neutron-rich radioactive nuclides have remained stubbornly inaccessible to precision study for decades because while fission of heavy elements abundantly produces these nuclides, their chemistry prohibitively impedes the release of all but the most volatile elements from traditional bulky fission sources [1]. The problem of chemistry dependence in neutron-rich beam production has been abated by the development of gas cells which can stop and thermalize reaction products in gas [2], with the most efficient being the new radio frequency (rf) gas catchers [3,4]. These developments have allowed comprehensive surveys of the properties of fission products to begin.

A key motivation for the study of these nuclides is the astrophysical rapid neutron-capture process (r process), which is thought to have produced half of the heavy nuclei in the universe [5–8]. The site of the r process which populates the universe is not known but is possibly within core-collapse supernovae or events which eject material from the crusts of neutron stars such as binary mergers. During an r -process event, the balance between neutron capture (n, γ) and photodissociation (γ, n) reactions determines the distribution of isotope populations within each element, while β decay moves nuclei to higher proton numbers. A critical nuclear physics input to the (γ, n) rate is the neutron separation energy (S_n) of the participant nucleus and is calculated from the atomic masses of parent and daughter isotopes. Direct mass measurements made with Penning traps [9] are only now reaching nuclides on possible r -process paths. The Canadian Penning Trap mass spectrometer (CPT) [10–12] and other Penning traps have made precision mass measurements of progressively more

neutron-rich nuclides [13–15] as the state of the art of gas-catcher technology has advanced.

Rf gas catchers are now sufficiently mature to justify new facilities dedicated to their beams. The new Californium Rare Isotope Breeder Upgrade (CARIBU) [16] is the first such facility, now online. CARIBU uses a ^{252}Cf spontaneous fission source in a large gas catcher to provide purified neutron-rich beams to either low-energy beam experiments or the Argonne Tandem-Linac Accelerator System (ATLAS) for reaccelerated beam experiments. This Letter describes the first science results using beams from CARIBU: mass measurements made with the CPT of the 33 nuclides shown in Fig. 1, 12 of which improve the precision over earlier attempts.

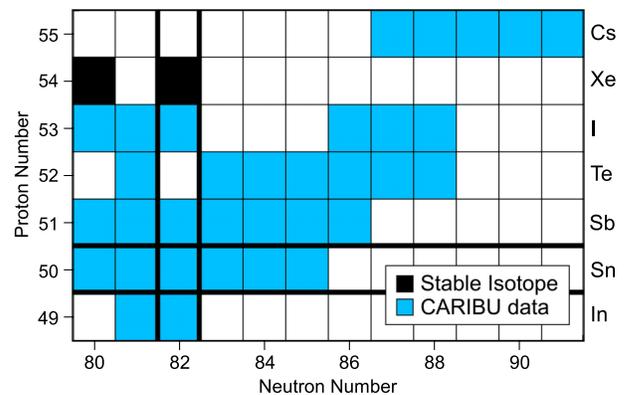


FIG. 1 (color online). Locations of the nuclides for which masses have been measured by the CPT at CARIBU. The shell closures at $N = 82$ and $Z = 50$ are highlighted.

The measurements.—For the bulk of these measurements, CARIBU used a 50 mCi fission source, providing beams to the CPT as intense as 4200 ions/s of ^{142}Cs , cleaned by an isobar separator [17] with a resolution of $\delta m/m \approx 1/9000$.

The CPT was moved to the CARIBU low-energy experimental area and recommissioned in 2011. Upgrades made during the move include the installation of a $^{133}\text{Cs}^+$ ion source for calibration and tuning, and liquid nitrogen cooling of the preparation linear radio frequency quadrupole ion trap located just before the CPT. Ions are ejected from the CARIBU buncher every 100–200 ms and are accumulated and cooled in the preparation trap until transfer to the CPT.

The CPT uses the time-of-flight ion cyclotron resonance method of Penning trap mass spectrometry for its measurements [18–20]. The cyclotron frequency $\omega_c = qB/m$ (where q is charge, B magnetic field strength, and m mass) is measured through successive attempts to resonantly convert ion bunches from slow to fast orbital motion. Ions are ejected out of the magnet to a detector, producing a time of flight versus applied frequency spectrum as shown in Fig. 2. The spectrum is fit with an approximation of the theoretical response in Ref. [21], with the minimum time of flight occurring at the cyclotron frequency. The mass is then determined from the ratio of the measured cyclotron frequency of a well-known calibrant species to that of the species of interest.

A strength of Penning trap measurements is their small systematic errors [19], which were investigated again for the recommissioning of the CPT. The superconducting magnet's field strength exhibited more scatter over time than before its move to CARIBU, such that the 11 ppb standard deviation among the calibrations was added as a systematic uncertainty to all measurements. Frequency shifts linear in mass difference are possible due to trap imperfections [9], and a search found such a shift in the

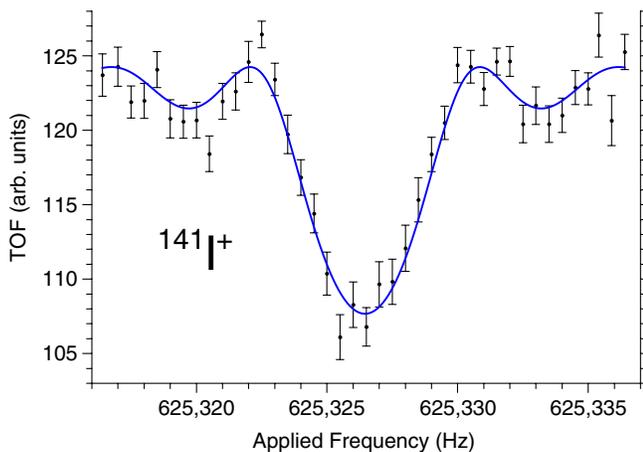


FIG. 2 (color online). Example data, a 200-ms excitation of $^{141}\text{I}^+$. The points represent the time-of-flight (TOF) data and the curve a fit to the theoretical response.

CPT, compensated for in the analysis by adjusting frequencies by $\Delta f/f = (A - 133) \times 1.74(18)$ ppb, where A is the mass number of the unknown. Finally, the detected ion rate was targeted below six detected ions per bunch in the trap to prevent ion-ion interactions [22] from introducing significant systematic effects.

Table I shows the masses measured by the CPT spectrometer. The known half-lives of these nuclides ranges down to 280 ms. For six of the nuclides— $^{130,131}\text{Sn}$, $^{132,134}\text{Sb}$, ^{133}Te , and ^{134}I —isomeric states were also clearly observed in the trap, which will be discussed in a future publication. For two isotopes $^{130,131}\text{In}$, isomeric states are expected to be present in the beam, but the CPT was unable either to resolve the states or to identify additional peaks, so no identification of the measured state or states has been made. Durations of the measurement excitations ranged from 0.1 to 7 s, depending on the lifetime of the nuclide and the proximity of isomers in cyclotron frequency. Total mass uncertainty ranges from 2.1 to 62 keV/ c^2 , with a median uncertainty of 8.6 keV/ c^2 .

In comparison with the 2003 Atomic Mass Evaluation (AME03) values [23], the previously seen trend of direct mass measurement values being higher than the AME03 far from stability continues (see examples in Refs. [10,13,15,32–34]). This is likely due to systematic flaws in the β end point measurements which previously dominated mass information in this region [35]. Also listed in Table I are comparisons to other direct measurements. The agreement for the seven nuclides measured by both the CPT and ISOLTRAP [25–28] is excellent. The 19 common ground-state measurements with JYFLTRAP [24] are mostly in agreement but with two masses—of ^{133}Te and ^{140}Te —disagreeing by 5.6σ and 3.2σ , respectively. Measurements of these two isotopes were repeated in separate runs at the CPT and gave consistent results. The known isomer in ^{133}Te is well resolved, and no known isomers are expected to interfere with the ground state measurements of either ^{133}Te or ^{140}Te . The cause of the disagreements remains unknown. Finally, the overlap with the nine ground states measured in the FRS-ESR ring [29,30] are in fair agreement, with two— ^{136}Te and ^{141}I —disagreeing by 2.0σ and 2.9σ , respectively.

Astrophysics impact.—The bulk of the neutron captures in the r process are thought to occur on time scales of order 1 s in environments of temperature $T \gtrsim 1$ GK and neutron density $n_n \gtrsim 10^{20} \text{ cm}^{-3}$ [7]. In that time, seed material from the nickel region must, through neutron capture and β decay, be processed up through the heavy elements to explain the observed elemental abundances. Because β decay is the process that moves material to higher proton numbers, the β -decay lifetimes of the nuclei involved determine how quickly the heaviest elements can be reached. The balance between neutron capture (n, γ) and photodissociation (γ, n) reactions determines the distribution of isotope populations within each element in a high-temperature r process.

TABLE I. CPT results from CARIBU beams. Mass values are combined with earlier CPT work where indicated, while all cyclotron frequency ratios are strictly from the new CARIBU-based results. Differences are shown from the AME03 [23] and other direct measurements from JYFLTRAP [24], ISOLTRAP [25–28], and the FRS-ESR ring [29,30]. In the difference columns, uncertainties in parentheses are the combined uncertainties, and those in curly brackets are of the sources cited. The 2003 AME is shown rather than the new 2012 AME [31] due to many of the latter’s entries being dominated by the measurements shown in columns 5–7 and in Ref. [15]. A # symbol indicates an extrapolated mass value in the AME03.

Nuclide	$\omega_c(^{133}\text{Cs}^+)/\omega_c(\text{Unknown}^+)$		Mass Excess of Neutral Atom (keV)						
	CPT at CARIBU	CPT	$\Delta_{\text{CPT-AME03}}$	$\Delta_{\text{CPT-JYFLTRAP}}$	Ref.	$\Delta_{\text{CPT-ISOLTRAP}}$	Ref.	$\Delta_{\text{CPT-FRS-ESR}}$	Ref.
¹³⁰ In ^a	0.977 576 25(16)	−69 652(20)	238(44){40}						
¹³¹ In ^a	0.977 575 78(16)	−67 876(35)	262(45){28}	149(35){2.6}	[24]				
¹³⁰ Sn ^b	0.977 491 606(29)	−80 130.8(3.6)	8(11){11}	2(5){4}	[24]	−3(16){16}	[27]		
¹³¹ Sn ^b	0.985 038 975(35)	−77 259.6(4.3)	55(22){21}	2(20){20}	[24]	4(11){10}	[28]	−22(120){120} ^a	[29]
¹³² Sn	0.992 568 893(22)	−76 549.0(2.8)	5(14){14}	−6(5){4}	[24]	−2(8){7}	[28]		
¹³³ Sn	1.000 138 949(29)	−70 869.1(3.6)	84(36){36}	5.3(4.3){2.4}	[24]	22(23){23}	[28]		
¹³⁴ Sn	1.007 698 87(13)	−66 444(16)	350(100){100}	12(16){4}	[24]	−120(150){150}	[28]		
¹³⁵ Sn	1.015 270 38(28)	−60 584(34)	210(400){400}#	48(35){3}	[24]				
¹³¹ Sb	0.985 000 799(84)	−81 986(10)	2(23){21}	−3(10){2.1}	[24]				
¹³² Sb ^b	0.992 543 975(49)	−79 633.8(6.1)	40(16){14}	1.8(6.7){2.7}	[24]			236(124){124} ^a	[29]
¹³³ Sb ^c	1.000 073 87(10)	−78 921.3(7.6)	21(27){25}	0(9){4}	[24]			−22(24){23}	[30]
¹³⁴ Sb ^{b,d}	1.007 637 742(82)	−74 012(10)	154(45){43}	10(10){2.1}	[24]				
¹³⁵ Sb	1.015 196 795(53)	−69 693.9(6.5)	14(100){100}	−4.3(7.1){2.9}	[24]			115(121){121}	[29]
¹³⁶ Sb	1.022 763 00(12)	−64 491(15)	390(300){300}#	19(16){7}	[24]				
¹³⁷ Sb	1.030 322 96(42)	−60 061(52)	200(400){400}#						
¹³³ Te ^b	1.000 041 770(52)	−82 899.8(6.5)	45(25){24}	38.4(6.8){2.2}	[24]				
¹³⁵ Te ^c	1.015 131 888(18)	−77 729.6(2.1)	98(90){90}	−1.7(3.3){2.6}	[24]			−5(123){123}	[29]
¹³⁶ Te ^c	1.022 682 783(51)	−74 423.3(3.7)	2(45){45}	2.4(4.7){2.9}	[24]			45(23){23}	[30]
¹³⁷ Te ^c	1.030 248 309(31)	−69 301.7(3.7)	260(120){120}	2.5(4.5){2.5}	[24]			−12(120){120}	[29]
¹³⁸ Te	1.037 801 624(61)	−65 695.3(7.6)	240(200){200}#	1(9){5}	[24]			60(122){122}	[29]
¹³⁹ Te	1.045 370 26(13)	−60 191(17)	610(400){400}#	14(17){4}	[24]				
¹⁴⁰ Te	1.052 923 64(50)	−56 577(62)	380(300){300}#	−220(68){27}	[24]				
¹³³ I ^e	1.000 017 873(52)	−85 858.2(6.4)	28(8){5}						
¹³⁴ I ^b	1.007 556 731(51)	−84 040.8(6.4)	32(10){8}						
¹³⁵ I ^e	1.015 083 028(17)	−83 778.9(2.0)	11(8){7}						
¹³⁹ I ^e	1.045 303 381(32)	−68 470.7(4.0)	367(31){31}					56(121){121}	[29]
¹⁴⁰ I	1.052 866 85(10)	−63 606(13)	670(200){200}#					−10(122){121}	[29]
¹⁴¹ I	1.060 420 75(13)	−59 927(16)	590(200){200}#					374(130){129}	[29]
¹⁴² Cs ^c	1.067 859 61(17)	−70 506.9(9.3)	8(14){11}			14(17){15} ^f	[25]		
¹⁴³ Cs	1.075 406 507(64)	−67 676.3(7.9)	−5(25){24}						
¹⁴⁴ Cs	1.082 966 39(25)	−63 256(31)	14(41){26}						
¹⁴⁵ Cs	1.090 516 40(13)	−60 057(16)	0(19){11}			−5(19){11}	[26]		
¹⁴⁶ Cs	1.098 078 820(69)	−55 323.2(8.6)	297(72){71}						

^a An unknown mixture of the ground and isomer states.

^b Resolved from the isomer; the ground-state mass is shown.

^c Mass is the combined result for the CPT measurement in Ref. [15] and the new CARIBU data.

^d The value in Ref. [15] has been superseded by CARIBU data due to resolution of the two states present.

^e The distant isomer was not targeted; the observed state is expected to be the ground state.

^f Result has been adjusted by us due to a change in the calibration value after Ref. [25]. See Ref. [15] for details.

In hotter environments, increased (γ, n) rates push this distribution closer to stability and thus to nuclides with longer β -decay lifetimes. At higher neutron densities, higher (n, γ) rates have the opposite effect.

Because of the scarcity of direct mass measurements on the r -process path, theoretical nuclear mass models are used to provide the S_n inputs needed for (γ, n) rate calculations. Unfortunately, even the most accurate mass models have rms mass errors near 500 keV/ c^2 for known nuclei [36–38], which, as is shown below, is insufficient for accurate r -process simulations. Some of the nuclides with masses

measured here—particularly isotopes of Sn and Sb—are on the r -process path for certain environmental conditions. For those paths, the measured Sn isotopes are among the nuclides with the longest β -decay lifetimes and are in a position to limit the progress of material to the heaviest elements. The waiting time at this critical point on the r -process path has never been calculated from directly measured accurate masses for realistic conditions.

We performed therefore a series of focused r -process simulations to determine the waiting times at Sn and Sb for a span of conditions, using both this new mass

information and, for comparison, three commonly used mass models: the finite range droplet model of 1995 (FRDM95) [39], the extended Thomas-Fermi plus Strutinsky integral with enhanced quenching model (ETFSI-Q) [40], and the recent Hartree-Fock-Bogoliubov-21 model (HFB-21) [37]. The code used is one modified from that used earlier by this group in a similar r -process calculation [41]. Within a given simulation run, the temperature and neutron density of the r -process environment are fixed. An initial population of nuclei is set on the low- N side at $N = 81$, and then the (n, γ) , (γ, n) , and β -decay processes are allowed to run over all appreciably populated isotopes of the simulated element. The (n, γ) and (γ, n) rates come from the mass-model-based calculations of Rauscher and Thielemann [42], which accept S_n values as input parameters for the (γ, n) rates. Because isotopes beyond those measured by the CPT participate to some extent for all simulated environmental conditions, AME03 masses were used for lower- N isotopes and mass model S_n values for higher N . To isolate the effect of the differences between each model and the experimental masses, simulations were run with each of the mass models appended to the CPT data. Separate simulations were run with each mass model used as an input. β -decay half-lives were taken from NuBase [43], where available, and model-based half-lives [44] were used for ^{138}Sn , ^{138}Sb , and beyond, scaled to match smoothly onto the known values. The simulations were run until half of the material had β decayed, providing a measure of the waiting time via this “effective” half-life in each condition. Simulations were run at 1.5 GK as a representative case and were repeated over a span of neutron densities until the new masses were no longer relevant.

Each mass-model-based simulation shows overly favorable prospects for moving past Sn, approaching a 1 s effective half-life at lower densities than the CPT-based simulations, as shown in the top of Fig. 3. This is due to the neutron separation energies of certain nuclei being too large in the mass models, impeding the (γ, n) reactions that would push material to longer-lived nuclei. In the ETFSI-Q case, this is caused mainly by a 350-keV error in the S_n of ^{133}Sn . The effect is more drastic for the FRDM95, which overestimates S_n of both $^{133,134}\text{Sn}$ by 260 and 635 keV, respectively, followed by an underestimation for ^{135}Sn by 340 keV. The combined effect of these is a prolonged span of densities for which ^{134}Sn is the dominant isotope, with its 1.1-s half-life setting the time scale. The effect in HFB-21 starts at higher density than in the other models, but is sharper. The delay is due to a small undershoot in the S_n of ^{133}Sn , and the rapid drop is started by a 604-keV error in ^{134}Sn . The mass models disagree with each other at ^{136}Sn over a span of 839 keV, which induces the differences among the CPT-based simulations apparent in the plot. Simulations at other temperatures showed the same effects but at higher density for higher temperatures, as expected.

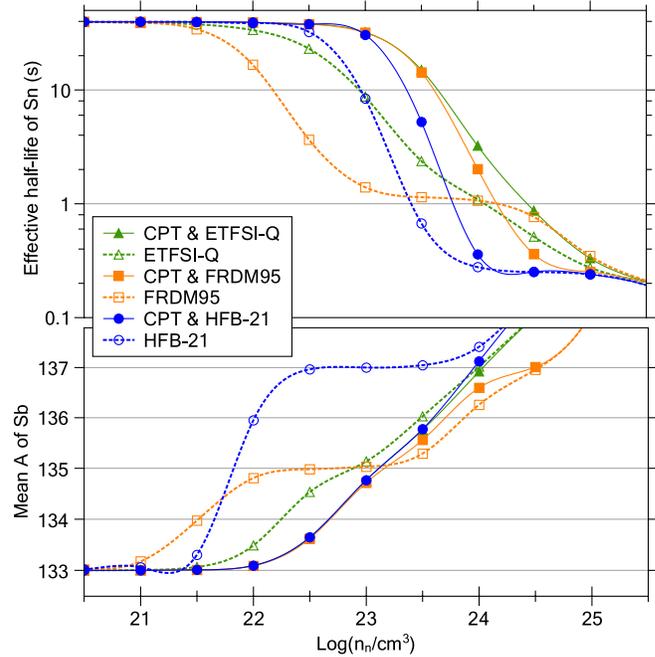


FIG. 3 (color online). Results of simulations performed at 1.5 GK. Top: The effective half-life of Sn versus neutron density. Each mass model tested predicts a drop in time at lower density than in reality. Bottom: The mean mass number of Sb versus neutron density. The mass models all predict a breakaway from the neutron shell closure earlier than in reality, with HFB-21 [37] showing an unnaturally wide expanse in which ^{137}Sb dominates. See the text for additional details.

Similar effects emerged from the simulations of Sb, but the effective half-life is nearly always shorter at Sn than Sb over the relevant conditions. Therefore, plotted in the bottom of Fig. 3 is the mean mass number of the Sb populations at the end of each simulation, to illustrate the impact on r -process abundances. There is a drastic effect in the HFB-21 simulations caused by an apparent lack of pairing around the neutron shell closure and consecutive 812 and 628 keV errors in the S_n of $^{136,137}\text{Sb}$, which almost instantly drives all material from ^{133}Sb to ^{137}Sb , an error of nearly four masses. Because this change is at low neutron densities, the primary effect will be during freeze-out, the time during which final r -process abundances (the only available r -process observable) are being determined.

These simulations provide a demonstration of the insufficiency of existing mass models for accurate r -process simulations. Progress in the development of mass models is slow, with the FRDM accuracy improving by 15% over the last 17 years [36] and the HFB models 32% over 9 years [37,45], for example. Therefore, without considerable improvements in nuclear theory, direct mass measurements may be the only sufficiently accurate source of S_n inputs to r -process simulations for the foreseeable future.

In an actual r -process event, the temperature and density will both be evolving rapidly as the material expands and reactions take place. These new results then provide a threshold of density which, for any given temperature, must be crossed in order to achieve a waiting time of ~ 1 s or less at Sn and move significant quantities of material to $Z > 50$ before the r -process event ends. Isotopes two neutrons farther from stability may be studied at CARIBU after installation of a 1 Ci fission source. This would extend Sn measurements as far as ^{137}Sn , which—with its 273-ms half-life—will set even stronger constraints on the r -process environment.

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