Overheating is a major bottleneck in optimizing the performance of microelectronic devices. Following Moore’s law, the power density in processors is increasing exponentially and calls for thermal management. With increased integration (current typical length scales being less than 45 nm base length), 3D packaging, and use of current controlled memory cells (examples of which include nonvolatile MRAM [1] or the “magnetic racetrack” [2]), these issues are coming to a head. Nevertheless, optimistic design engineers promise batteries that last for 24 hours in the next generation of laptop computers, with minimal heat dissipation at highest performance. What strategies do they have in mind?

One promising route to this goal is purely electric-field-controlled devices, which cause much less Joule heating than magnetic devices that are switched with a current. This technological need revived interest in magnetoelectric (ME) materials [3], in which magnetism can be controlled via electric fields, and vice versa [4]. Its promise reaches its zenith in multiferroics. These materials that undergo phase transitions into long range magnetic and electric order are potentially the most promising materials for electric control of magnetism, for storage and for processing. What strategies do they have in mind?

Unfortunately, the linear magnetoelectric effect is so small in most materials [3] that this direct approach has no chance of being successful in practice. However, much effort has been put into related devices. Binek and Doudin [6] were probably the first to draw attention to the potential of spintronic concepts in this pursuit, when they proposed to use magnetoelectric materials as active tunnel barriers or pinning layers of spin-dependent transport devices like magnetic random access memory (MRAM) [1]. In contrast to Fig.1, $\Delta m_{\text{ME}}$ now controls the magnetization state of adjacent ferromagnetic (FM) layers via exchange coupling, which is the crucial keyword in all actual device proposals. It is this that has motivated a paper published in Physical Review Letters by Lebeugle et al.[7], who have demonstrated electric field switching of the magnetic anisotropy of a soft magnetic layer of Ni$_{0.78}$Fe$_{0.22}$ (“NiFe” or Permalloy) attached to a single crystal of ferroelectric and antiferromagnetic BiFeO$_3$ (BFO).

The results by Lebeugle et al. should be regarded in parallel to very similar experiments of Ramesh’s group [8] on soft magnetic layers of Co$_{0.9}$Fe$_{0.1}$ (“CoFe”) attached to thin films of BFO. Both composites involve the canonical multiferroic BFO [9], still the only room-temperature multiferroic with exceptionally high ordering temperatures: ferroelectric Curie temperature $T_C = 1103$ K and antiferromagnetic Néel temperature $T_N = 643$ K.

In both Refs. [7] and [8], it has been shown that an electric field-induced change of the ferroelectric polarization of the BFO substrate is able to toggle the easy direction of the magnetization in the ferromagnetic layer by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited. The first is the magnetoelectric coupling within BFO by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited. The first is the magnetoelectric coupling within BFO by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited. The first is the magnetoelectric coupling within BFO by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited. The first is the magnetoelectric coupling within BFO by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited. The first is the magnetoelectric coupling within BFO by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited. The first is the magnetoelectric coupling within BFO by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited. The first is the magnetoelectric coupling within BFO by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited. The first is the magnetoelectric coupling within BFO by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited. The first is the magnetoelectric coupling within BFO by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited. The first is the magnetoelectric coupling within BFO by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited. The first is the magnetoelectric coupling within BFO by use of the magnetoelectric effect. More precisely, two successive coupling mechanisms are exploited.
interactions at the interface between the antiferromagnetic BFO and the ferromagnetic Permalloy (or CoFe). More precisely, it is the projection of the antiferromagnetic order that couples to the ferromagnetic magnetization. Slightly different interpretations are in order, since the Fe$^{3+}$ spins of BFO form cycloids in the single crystal [7], while they are uniformly canted in the thin film [8]. The cycloidal propagation vector favors a parallel alignment with the Permalloy spins. As a result the effective coupling is several times larger in the crystal than in the film. Unfortunately, the magnetization of the Permalloy layer could not be completely switched along the directions of the anisotropy axes, as it is impossible to form a ferroelectric single domain in the (001) plane of BFO. Out of the manifold of eight differently poled domains, only four of them can be selected by a uniform inplane field. Similar results emerged from the switching experiments on BFO/CoFe bilayers [8].

At this point one might ask whether BFO really is the most promising material for switching magnetism with an electric field. The remaining ambiguity of electric switching due to crystallographic multiaxiality obviously calls for a uniaxial magnetoelectric material exhibiting only two different 180° ferroelectric, ferromagnetic, or antiferromagnetic domains. We found that the primordial magnetoelectric material, chromium oxide Cr$_2$O$_3$[12], fulfills this requirement in an unprecedented way. Its two uniaxial antiferromagnetic domains can be switched magnetoelectrically, by simultaneously applying a magnetic and an electric field, $\mathbf{E}$ and $\mathbf{H}$, on cooling to below $T_N = 308 \text{ K}$[13] or isothermally in a large enough field product $E\cdot H$[14]. By keeping $H$ constant (e.g., with a permanent magnet) the antiferromagnetic single domain state can be switched with the electric field $\pm E$. This has spectacular consequences to an attached ferromagnetic layer, whose magnetization is also switched by 180° due to the so-called exchange bias effect [15]. We showed this on a ferromagnetic trilayer Pt/Co/Pt on top of a (0001) oriented single crystal of Cr$_2$O$_3$[16]. Following earlier ideas [6], the design of magnetoelectric random access memory and logic cells, MERAM and MEXOR [17], was obvious. Their binary information storage is based on the low and high tunnel magnetoresistances of the magnetoelectric switchable “free layer” coupled by a tunnel barrier to a “pinned layer” with fixed ferromagnetic magnetization, parallel or antiparallel to the “free” one, respectively. The next steps to take toward possible applications will be the reduction of leakage currents in Cr$_2$O$_3$ thin films and the increase of their Néel temperature, e.g., by alloying with Fe$_2$O$_3$.

In conclusion, the present strategies of switching magnetism with electric fields—targeting either ferromagnetic anisotropy or exchange bias—while still in their infancy, are highly promising, though both fundamental and material issues still have to be solved.

**References**


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Wolfgang Kleemann is professor of experimental and applied physics at the University of Duisburg-Essen. He received his Ph.D. in 1968 from the University of Göttingen and became professor at the Universities of Paderborn (1975) and Duisburg (since 1982). His perpetual research interest in both magnetism (ferro- and antiferromagnetic crystals, multilayers and nanoparticles) and ferroelectricity (crystals, films and ceramics) brought him to Universities like Paris-Sud, UC Santa Barbara, RIKEN Tokyo, SNU Seoul, and UFMG Belo Horizonte as a visiting professor. His activities have recently become focused on magneto-electric and multiferroic materials.