



Entanglement-Enhanced Classical Communication Over a Noisy Classical Channel

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We present and experimentally demonstrate a communication protocol that employs shared entanglement to reduce errors when sending a bit over a particular noisy classical channel. Specifically, it is shown that given a single use of this channel, one can transmit a bit with higher success probability when the sender and receiver share entanglement compared to the best possible strategy when they do not. The experiment is realized using polarization-entangled photon pairs, whose quantum correlations play a critical role in both the encoding and decoding of the classical message. Experimentally, we find that a bit can be successfully transmitted with probability 0.891 ± 0.002 , which is close to the theoretical maximum of $(2 + 2^{-1/2})/3 \approx 0.902$ and is significantly above the optimal classical strategy, which yields $5/6 \approx 0.833$.

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Two parties that share an entangled quantum system can achieve communication tasks which would otherwise be impossible: Sending two bits of classical information using only one qubit [1], unconditionally secure communication [2], transferring quantum information from one quantum system to another using two classical bits [3], and reducing communication complexity in distributed computations [4].

It is much less studied how entanglement can assist in sending classical information over a classical channel. The main result in this context is that entanglement cannot increase the capacity of any classical channel [3,5] in the sense of Shannon [6]. However, it has recently been shown that the number of different messages that can be sent error-free with a single use of a noisy classical channel (i.e. the *one-shot zero-error capacity*) can be increased when the sender and receiver share entanglement [7]. Another interpretation of this result is that for a fixed number of possible messages and a fixed number of channel uses, entanglement can be used to increase the probability of successful decoding, which raises interesting general questions about when entanglement can assist in this way.

In this Letter we describe an example of such an entanglement-enhanced classical communication protocol that, compared to the example given in [7], exhibits a larger absolute gap in the assisted and unassisted success probabilities, involves a much simpler classical channel, a smaller entangled state and, most significantly, is experimentally feasible. We go on to describe an experimental implementation which clearly demonstrates the advantage gained from using entanglement. This represents an application of entanglement in a new setting, where it is surprising that one can benefit from quantum effects.

We first briefly discuss error-correcting codes for classical channels with and without entanglement assistance. For a specific classical channel, which we label $\mathcal{N}_{\boxtimes}$, we show that the maximum success probability for sending 1 bit with a single use of this channel is $5/6 \approx 0.833$, even

if the parties have shared randomness at their disposal. By employing a pair of maximally entangled qubits, this probability can be increased to $(2 + 2^{-1/2})/3 \approx 0.902$. Using entangled photons we then implement this protocol experimentally, achieving a success probability of 0.891 ± 0.002 . This is close to the theoretical value and significantly better than the best classical strategy.

Classical vs entanglement-assisted codes.—The task we are studying is this: Alice wants to be able to send one of M possible messages to Bob, by making one use of a noisy classical channel $\mathcal{N}$ (and no other signals). The channel has a finite set of input symbols, and a finite set of output symbols, and is described by the conditional probabilities of the output symbols given the input symbol. Given a uniform prior distribution on the messages, we want to maximize the probability that Bob determines the message correctly.

If we have a purely classical protocol which depends on some random variable R , its success probability will be the mean of the success probabilities of the protocols for fixed values of R . Therefore, it is no larger than the best success probability attained by some fixed value of R and we can eliminate the dependence on R without detriment (note that this is true even if R is shared randomness, i.e., its value is known to both parties). Therefore, there is an optimal classical protocol which is deterministic: Alice maps each message to some input symbol, and Bob has a decoding rule that maps output symbols to messages (if the channel is of the form $\mathcal{N}^{\otimes k}$, corresponding to k independent uses of a given channel $\mathcal{N}$, then such a protocol is called a *block code* of block length k for the channel $\mathcal{N}$).

On the other hand, we will see that shared entanglement can increase the probability of success beyond any classical protocol. The most general form of an entanglement-assisted protocol to send one of M messages using a classical channel is illustrated in Fig 1: Alice and Bob possess systems A and B in some entangled state ρ_{AB} (note that a separable state can be simulated by shared randomness). The message q determines which

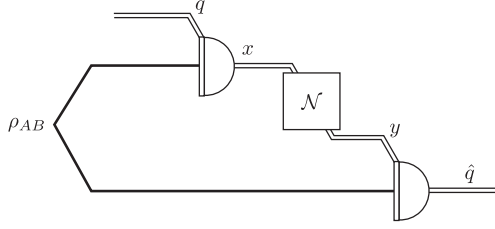


FIG. 1. General form of an entanglement-assisted coding protocol for sending classical message q with a single use of a classical channel $\mathcal{N}$ and an entangled state ρ_{AB} . Double lines represent classical communication.

measurement Alice performs on her system A and the measurement outcome determines the input x she makes to the channel. This does not imply that the channel input can not also depend on q , but this dependence can be absorbed into the definition of the measurement. The channel output y then determines the measurement Bob does on his system, B , and its outcome determines his decoding $\hat{q}$ of the message (again, the definition of the measurement can incorporate any dependence on y). For more details, see [7].

The channel.—The particular noisy classical channel $\mathcal{N}_{\boxtimes}$ we consider in this Letter takes a two-bit input, (b_1, b_2) . It produces an output (t, b) consisting of a trit t which is selected at random from the set $\{1, 2, P\}$ (each with probability $1/3$), and a bit b which is equal to b_1 when $t = 1$, b_2 when $t = 2$, and is the parity $b_1 \oplus b_2$ when $t = P$. The conditional probability matrix $\Pr[(t, b)|(b_1, b_2)]$ of $\mathcal{N}_{\boxtimes}$ is shown in Fig. 2(a). In Fig. 2(b), the four input symbols are represented as the circular vertices of a graph whose edges (with square labels) correspond to the six output symbols. In terms of this picture, when an input is made, the output is chosen uniformly at random from the three edges in the graph which connect to it.

A classical code to send a bit with one use of $\mathcal{N}_{\boxtimes}$ uses two of the four possible inputs to represent the two messages “0” and “1”. For example, we might use $(0, 1)$ to encode “0” and $(1, 0)$ for “1”. The optimal decoding map is easy to see: If the output is $(1, 0)$ or $(2, 1)$ then the input was certainly $(0, 1)$ so Bob decodes “0”. Likewise, he decodes $(1, 1)$ and $(2, 0)$ to message “1”, and this is also always correct. Output $(P, 0)$ never occurs. If the output is $(P, 1)$ then the two messages are equally probable, so the

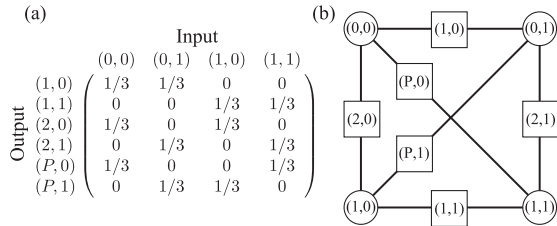


FIG. 2. The conditional probability matrix of $\mathcal{N}_{\boxtimes}$ (a) and a graph representing its structure (b). The inputs (outputs) are labeled as circles (squares).

best Bob can do is guess. He will be right with probability $1/2$, and this output occurs with probability $1/3$ so he is wrong with probability $1/6$. The symmetry of the channel [which is evident in Figure 2(b)] means that any other assignment (or code) is equivalent to this one under relabeling of the channel outputs, and has the same probability of successful decoding, namely, $5/6 \approx 0.833$.

The entanglement-assisted protocol.—Shared entanglement allows us to improve on the optimal classical strategy described in the last section. The protocol assumes that the two parties each have one photon from a maximally polarization-entangled state of two photons $|\Phi^+\rangle = (|HH\rangle + |VV\rangle)/\sqrt{2}$. For the protocol to proceed, each party performs a measurement in one out of two bases on their photon. Each measurement has two possible outcomes, projecting on quantum states $|\theta\rangle = \cos\theta|H\rangle + \sin\theta|V\rangle$. We will represent each such measurement by $(|\theta_0\rangle, |\theta_1\rangle)$ to denote the (orthogonal) outcomes 0 and 1 for each measurement, respectively.

For Alice, measurement 0 is $(|\pi/4\rangle, |3\pi/4\rangle)$ and measurement 1 is $(|0\rangle, |\pi/2\rangle)$. Bob’s measurement 0 is $(|\pi/8\rangle, |5\pi/8\rangle)$ and measurement 1 is $(|3\pi/8\rangle, |7\pi/8\rangle)$. Note that these are the measurement settings that give rise to a maximal violation of a Clauser-Horne-Shimony-Holt Bell’s inequality [8]. Alice performs measurement $q = 0$ or 1 and obtains outcome α while Bob makes measurement $v = 0$ or 1 and obtains outcome β . Here α and β are bit values corresponding to the measurement outcomes. The marginal distributions of α and β are uniform, independent of the inputs, and the correlations will be such that the relation [9]

$$\alpha \oplus \beta = qv \quad (2)$$

holds with probability $\omega = (1 + 2^{-1/2})/2$ [10]. To send the bit q , Alice performs measurement q , randomly

TABLE I. The rows of this table track the protocol from left to right for the three possible values of t : Bob obtains (t, b) from the channel, b ’s dependence on the message q and Alice’s measurement outcome α being fixed by t as described in the channel section; Based on t , Bob determined his measurement choice v . He gets the outcome β —in this column we assume that the relation (2) does hold. Substituting the choice Bob makes for v in this relation determines β as a function of α and q . The next column shows the function of b and β Bob calculates to get his decoding of the message $\hat{q}$. When the relation (2) holds, this function is equal to q for all t . If the relation does not hold, then if $t = 1$ (with probability $1/3$) he is still correct. Therefore the probability of successful decoding is $1/3 + 2\omega/3 = (2 + 2^{-1/2})/3$. The last two columns show the settings of the Pockels cells used to implement Bob’s measurement in the experiment (1 is “ON”, 0 is “OFF”).

t	b	Bob Chooses $v =$	β	$\hat{q}$	X_{on}	Z_{on}
1	q	Irrelevant	n/a	b	n/a	n/a
2	α	1	$q \oplus \alpha$	$b \oplus \beta$	$1 \oplus b$	b
P	$q \oplus \alpha$	0	α	$b \oplus \beta$	b	b

obtaining some outcome α . She then uses (q, α) as the input to $\mathcal{N}_{\boxtimes}$. In Table I and its caption we describe Bob's part of the protocol and how the success probability of $(2 + 2^{-1/2})/3 \approx 0.902$ can be achieved [11].

Experimental realization.—The setup of our experiment is shown in Fig. 3. We implement the classical channel, $\mathcal{N}_{\boxtimes}$, with its four inputs and six outputs by using an electronic circuit (CPLD logic chip Xilinx Xc2c64a) that is controlled by a random trit, t , is generated at 200 kHz by LabView's pseudorandom number generator (RNG). We generate the entangled resource state for our communication protocol by using type-II spontaneous parametric down-conversion. A 0.7 mW diode laser at 404 nm pumps a 25 mm periodically-poled KTiOPO₄ (PPKTP) crystal in a Sagnac configuration, emitting entangled photons which are subsequently single-mode fiber coupled [12,13]. Typically we observe a coincidence rate of 15 kHz directly at the source.

Alice's message, q , is selected randomly by means of a beam splitter; if her photon is transmitted she performs measurement 0, and if the photon is reflected she performs measurement 1. The input to $\mathcal{N}_{\boxtimes}$ is determined by which of the detectors fires, as described in the last section and illustrated in Fig. 3. Depending on the output y of the channel, Bob needs to actively choose his measurement basis and may need to invert his measurement outcome in order to decode the message. These actions are implemented using two fast RbTiOPO₄ Pockels cells, aligned for σ_x (X) and σ_z (Z) operation, respectively [13,14]. The states X_{on} and Z_{on} of the Pockels cells, which can each independently be 0 or 1, are shown in Table I.

After passing the Pockels cells, Bob's photon is measured in the $(|\pi/8\rangle, |5\pi/8\rangle)$ basis, where the output β is 0 or 1 depending on whether the photon is detected in the transmitted or reflected output port of the analyzer module,

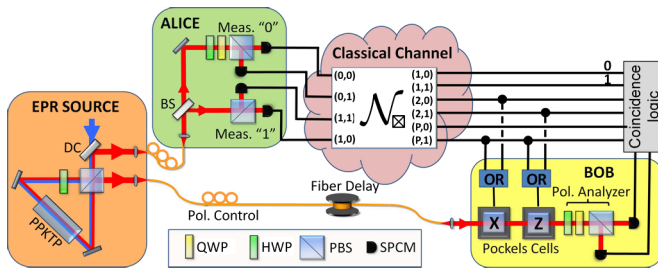


FIG. 3 (color online). Experimental implementation. The Sagnac-spontaneous parametric down-conversion source generates maximally entangled photon pairs of the form $|\Phi^+\rangle$, which are coupled into single-mode fibers (SMF). One of the photons is brought to Alice's side, where two polarization analyzer modules ($|H/V\rangle$ and $|+/-\rangle$ basis, respectively) are separated by a 50/50 beam splitter (BS). The classical channel $\mathcal{N}_{\boxtimes}$ performs its mapping dependent on the RNG operated at 200 kHz. This ensures that every new input to the channel will likely encounter a different random signal. Depending on the output of $\mathcal{N}_{\boxtimes}$, fast Pockels cells perform σ_x (X) or σ_z (Z) on Bob's photon, which is delayed in a 50 m SMF to account for the feed-forward time.

respectively. Unless $t = 1$, Bob uses the outcome of the measurement in combination with the output b of the channel to decode the message. If $t = 1$, Bob ignores the measurement result and directly uses b as the decoded message $\hat{q}$ (see Table I). The measurement bases of Alice and Bob are set by half- and quarter-wave plates (HWPs, QWPs) followed by polarizing beam splitters (PBSs).

To evaluate the quality of our entangled state, we performed quantum state tomography [15,16]. On Alice's side, we analyzed the photons in the transmitted arm of the BS. We recorded coincidences between this output and Bob's polarization analyzer following the (switched off) Pockels cells. Coincidence measurements were integrated over 8 s for each of 36 different measurements, comprising all combinations of the six eigenstates of X , Y , and Z on Alice's and Bob's qubit, respectively. Using a maximum-likelihood technique [16] we reconstructed the density matrix shown in Fig. 4.

To characterize the performance of the channel during the experiment, we record the coincidence events for all possible input-output combinations. The frequencies of the outputs for each of the inputs are shown in a "truth table" in Fig. 4. We can quantify the overlap between the measured truth table, N_{exp} and the ideal truth table, N_{th} , using the so-called inquisition [17], $I = \text{Tr}(N_{\text{exp}} N_{\text{th}}^T) / \text{Tr}(N_{\text{th}} N_{\text{th}}^T)$. For our channel, $I = 0.9992 \pm 0.0001$, where the uncertainty was calculated by a Monte Carlo simulation with binomially distributed random-signal frequencies (which stem from the RNG in the experiment) added in each run.

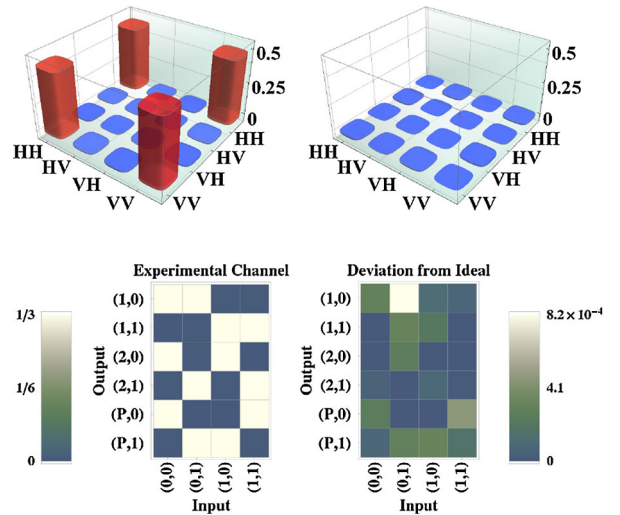


FIG. 4 (color online). Source and channel characterization. Top: Experimentally reconstructed density matrix of our two-photon entangled state ρ_{exp} : real part (left panel) and imaginary part (right panel). It has fidelity $F = 0.981 \pm 0.001$ with the ideal $|\Phi^+\rangle$ and a tangle of $\tau = 0.925 \pm 0.004$. The error bars for these results were calculated by a 200 run Monte Carlo simulation, adding Poissonian noise to the count statistics in each run. Bottom: Truth tables for the experimental channel (left panel) and its deviation from an ideal one [right panel—c.f. Fig. 2(a)], as recorded during our experiment.

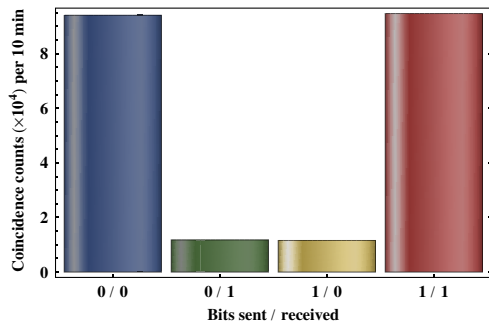


FIG. 5 (color online). Experimental results. The bar chart displays coincidences recorded (input bit, decoded bit) over a period of 10 minutes. In total, we recorded 188 845 (23 429) bits that were successfully (unsuccessfully) transmitted. These data yield a success probability $P_{\text{exp}} = 0.891 \pm 0.002$.

Experimental results.—In our experiment, we record all combinations of coincidence counts between Alice’s four single-photon counting detectors and the two on Bob’s side. This allows us to obtain the success probability P_{exp} , i.e., the ratio of successfully received and decoded bits over the total number of bits sent. The results are shown in Fig. 5. From the data we obtain $P_{\text{exp}} = 0.891 \pm 0.002$, where the error bar stems from Poissonian count statistics. Imperfect state creation and feed-forward operations lead to a decrease from the ideal, theoretical success probability, $P_{\text{th}} = (2 + 2^{-1/2})/3 \approx 0.902$. The tomographic data of our entangled resource state allow us to calculate the expected success probability of our protocol, which we infer as 0.8957 ± 0.0004 . This shows that our implementation of the protocol is mostly limited by the quality of our entangled resource state.

Conclusion.—In summary, we have shown how shared entanglement can be used to improve the performance of a completely classical communication task, namely, sending one bit with a single use of a noisy classical channel. This scheme shows how entanglement can offer a distinct advantage in a classical error coding scenario involving a finite number of channel uses. Our results lead to interesting questions for further study: which classical communication channels can benefit from entanglement and by how much? For instance, can we find general bounds on the gap between the error probability with and without entanglement assistance? How do these ideas generalize in the

context of multiterminal communication assisted by multipartite entanglement?

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